

Breaking the mass-sheet degeneracy with gravitational wave interference in lensed events

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1. Background

The MSD is a well-known problem in lensing of light.

It consists on [1]

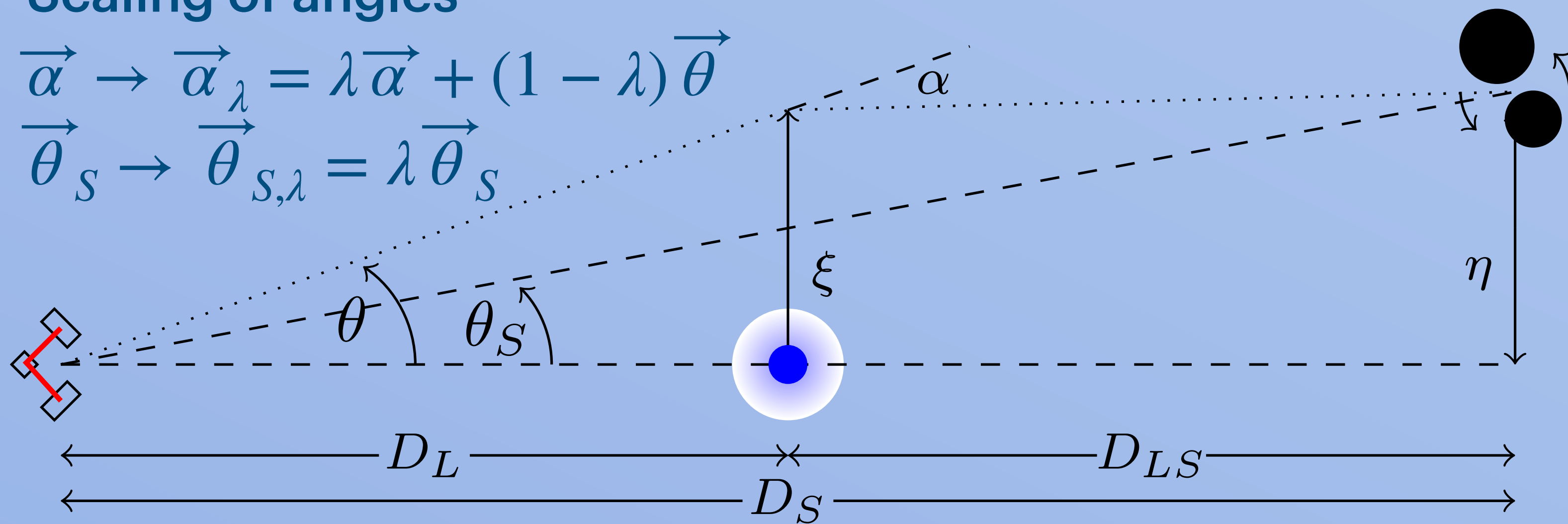
- Scaling of lens mass

(κ - surface mass density)
 $\kappa \rightarrow \kappa_\lambda = \lambda\kappa + (1 - \lambda)$

- Scaling of angles

$$\vec{\alpha} \rightarrow \vec{\alpha}_\lambda = \lambda\vec{\alpha} + (1 - \lambda)\vec{\theta}$$

$$\vec{\theta}_S \rightarrow \vec{\theta}_{S,\lambda} = \lambda\vec{\theta}_S$$



2. Problem

- Observables are preserved!
- Biased estimations of lens parameters
- Biased estimation of cosmological parameter (e.g. H_0)

3. Solution

In lensing of light: multiple images; independent mass estimation of lens; multiple lenses.

In GW lensing: 1 image and 1 lens can break MSD!

4. Gravitational Wave lensing

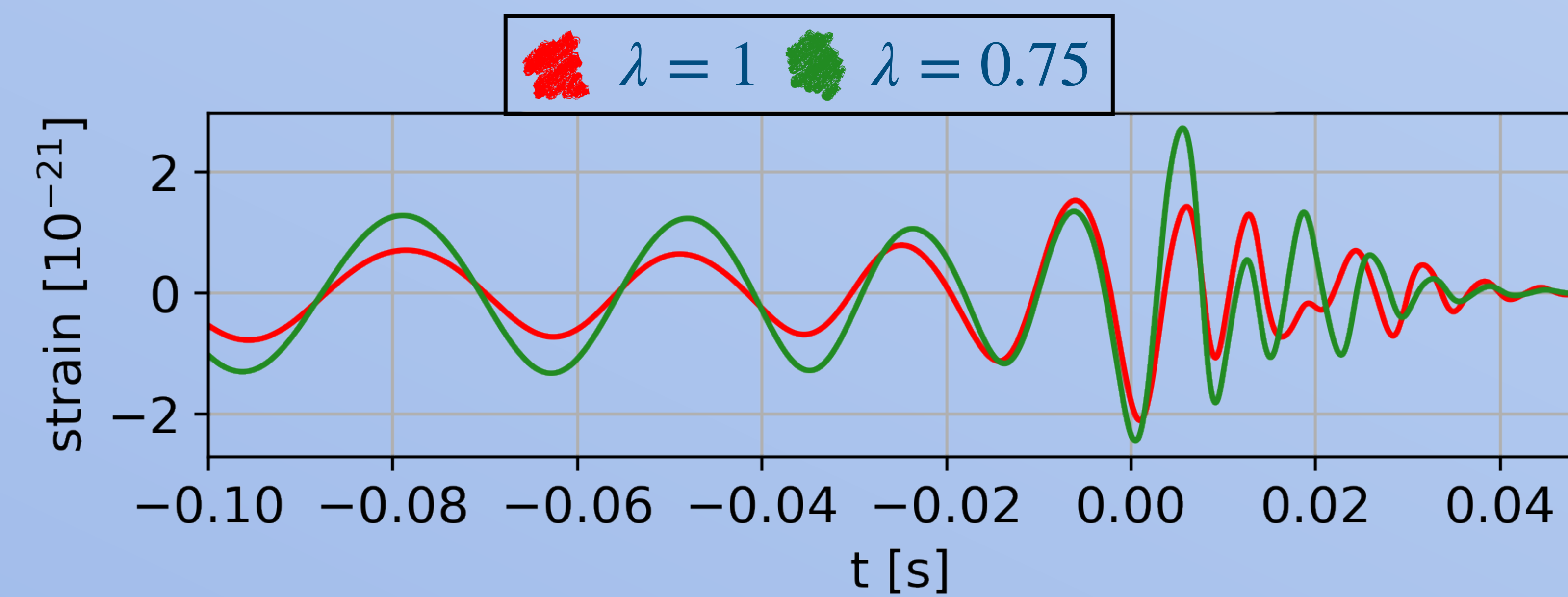
Lensed waveform in frequency domain $\tilde{h}_L(f)$ is:

$\tilde{h}_L(f) = \tilde{h}(f) \cdot F(f, \theta_s)$, with $\tilde{h}(f)$ the unlensed signal, f the GW frequency and $F(f, \theta_s)$ the amplification factor [2], which contains all info about lens.

5. MSD in GW lensing

According to the MSD transformation, also the amplification factor changes: $F \rightarrow F_\lambda$.

Signal with $M_S = 100 M_\odot$ at $z_S = 0.1$ lensed by a point mass lens with $M_L = 500 M_\odot$ & $y (\sim \theta_S) = 1$ at $z_L = 0.01$, gives



The shape of the lensed waveform depends on λ . Hence, the degeneracy is broken!

We studied, then, up to which level the MSD is broken.

6. Template matching (signal-to-noise ratio - SNR)

In a matched filtering analysis, the SNR is

$$\rho \approx \frac{(h | h_T)}{\sqrt{(h_T | h_T)}}$$

with h the signal,

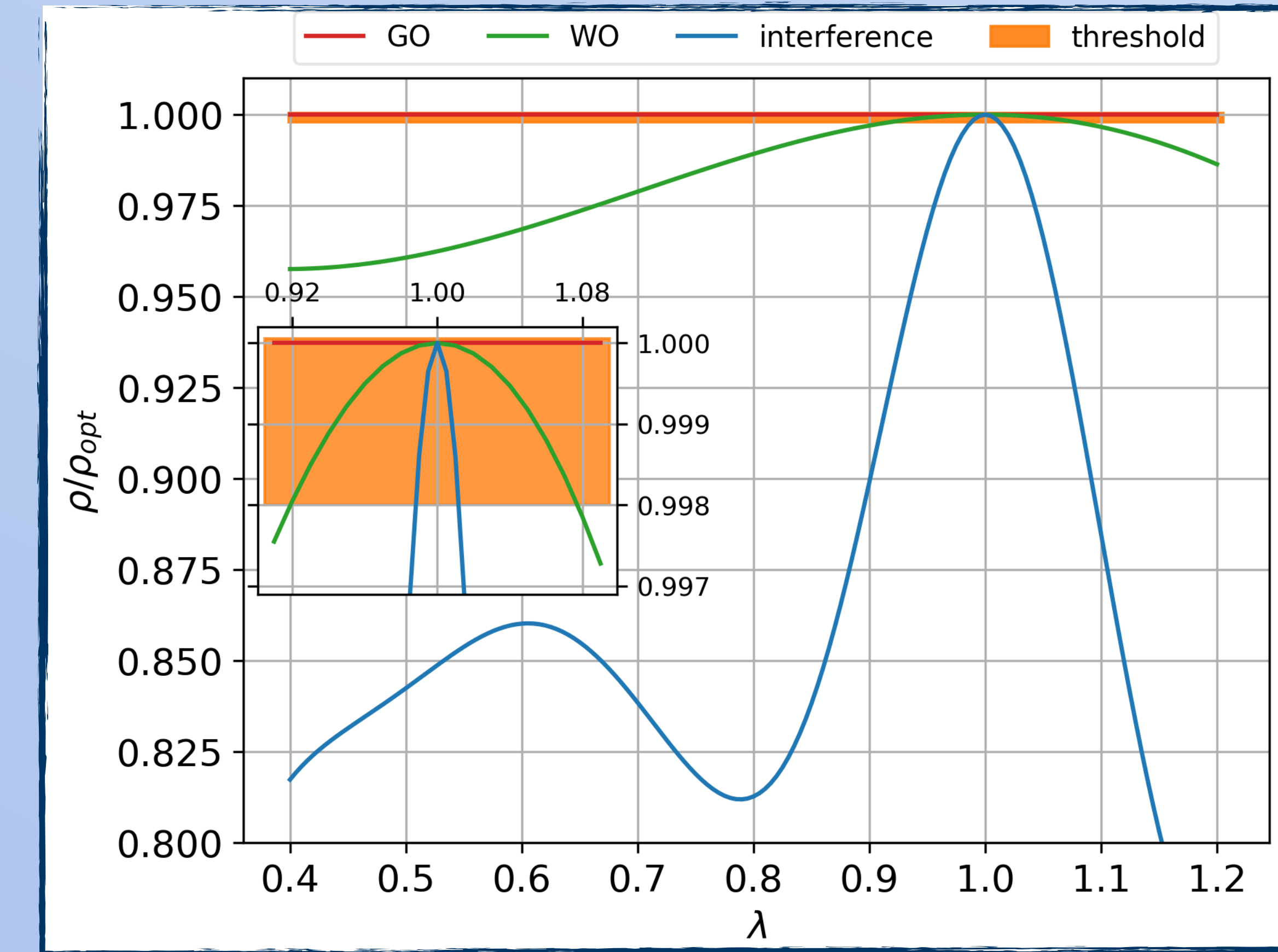
h_T the template and $(a | b)$ the inner product.

$$\frac{\rho}{\rho_{opt}} = 1 - \left[\frac{1}{2} \frac{\Delta\chi^2}{\rho_{opt}^2} \right],$$

with $\Delta\chi^2 \approx 11.8$ at 3σ . Then, we compare the signal with the (supposedly) degenerate one, computing ρ/ρ_{opt} .

$$(a | b) = 4 \text{Re} \left[\int_0^\infty \frac{\tilde{a}(f) \cdot \tilde{b}^*(f)}{S_n(f)} df \right]$$

with $S_n(f)$ - (single-sided) power spectral density (L1-O3-LIGO)



In the Figure, we show the behaviour of the match between the original signal and the degenerate one (ρ/ρ_{opt}) with λ .

When the curve leaves the confidence region, the MSD is broken, like in interference and wave-optics regimes (blue and green curves), but not in geometrical optics (red curve).

7. Results

The errors on the inferred parameters of the lens are:

- For a signal with $\rho = 11$, $\Delta y < 40\%$ and $\Delta M \approx 35\%$
- **For a signal with $\rho = 55$, $\Delta y \approx 5\%$ and $\Delta M \approx 6\%$**
- From dynamics calculations [4]: $\Delta M \approx 12 - 20\%$

8. Conclusions

- MSD ruins inference precision in lensing events
- In some cases of GW lensing (interference and wave-optics regimes), the MSD can be broken
- How well it is broken depends on the strength of the signal and sensitivity of detectors.

Bibliography

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- [3] M. Maggiore, "Gravitational Waves: Volume 1" (OUP Oxford, 2008)
- [4] P. Schneider and D. Sluse, Astron. Astrophys. 559, A37(2013)

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